

# Lifetime of breathers in hard and soft potentials

Juan FR. Archilla<sup>†</sup>, Jānis Bajārs<sup>‡</sup>, Sergej Flach<sup>§</sup>,

<sup>†</sup> Group of Nonlinear Physics, Department of Applied Physics I  
Universidad de Sevilla, ETSII, Avda Reina Mercedes s/n, 41012-Sevilla, Spain

<sup>‡</sup> Faculty of Physics, Mathematics and Optometry  
University of Latvia, Jelgavas Street 3, Riga, LV-1004, Latvia

<sup>§</sup> Center for Theoretical Physics of Complex Systems, Institute of Basic Science  
Expo-ro 55 Yuseong-gu, Daejeon 34126, South Korea  
Email: archilla@us.es, janis.bajars@lu.lv, sflach@ibs.re.kr

**Abstract**—We study the lifetime of breathers when they are produced in a thermalized background and the evolution to thermalization is observed. We deduce a theoretical value of a magnitude called the participation number  $P$  that indicates the arrival to thermal equilibrium. It is found that the lifetime has a great variability, but their average have an exponential dependence of the breather energy. This property holds for several systems although the exponential coefficients differ from system to system.

## 1. Procedure and results

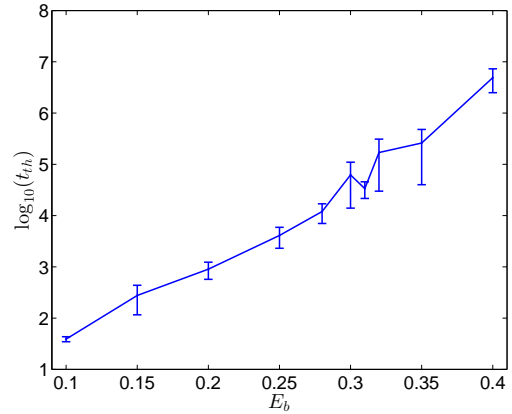
We consider both symmetric hard and soft Klein-Gordon potentials and harmonic coupling. The Hamiltonian of the system in scaled variables is given by:

$$H = \sum_n \frac{p_n^2}{2} + \omega_0^2 \left( \frac{u_n^2}{2} + s \frac{u_n^4}{4} \right) + \varepsilon \frac{1}{2} (u_{n+1} - u_n)^2,$$

where  $\omega_0$  is the frequency of the isolated oscillator at the linear limit and it is taken as  $\omega_0 = 1$ , and  $\varepsilon = 0.05$  is the coupling parameter. The nonlinearity parameter  $s = \pm 1$  selects a hard or soft potential. Other values of the parameters can be scaled to the previous equation, except for the ratio  $\varepsilon/\omega_0^2$ .

We first obtain breathers [1] from the anticontinuous limit [2] to have an indication of the energies involved. We can deduce the frequencies and their stability, being  $\omega_b \in (1.1, 2)$  for hard potentials and  $\omega_b \in (0.54, 1)$  for soft potentials. The energies are of a few tenths depending of the frequencies but smaller for soft potentials.

A measure of localization is the participation number  $P = E^2 / \sum e_n^2$ , where  $e_n$  is the local energy and  $E = \sum e_n$ .  $P$  takes values between 1, when all the energy is only at one particle and the number of particles  $N$  if all of them have the same energies. We deduce that at thermal equilibrium



$P = N/2$ . Then we thermalize the system with a given average local energy  $\langle e_n \rangle = 0.02$  and add a kinetic energy  $E_b$  at a single site, thereafter observing the time until  $P$  becomes  $N/2$ . These times vary greatly but their averages follow an exponential law as can be seen in the figure. This confirms that K-G systems have similar relaxation to equilibrium as the DNLS in spite of not having a conserved norm [3].

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ORCID iDs JFR Archilla: 0000-0001-6583-6114,

J Bajārs: 0000-0001-7601-8694,

S Flach: 0000-0003-1710-3746



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